

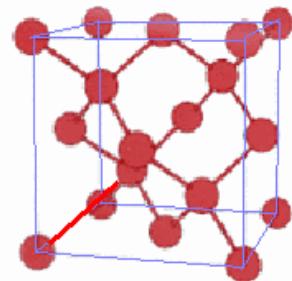
C_{60} , 碳纳米管, 石墨烯 物性与器件研究

物理所 吕 力

2009年12月17日 北大

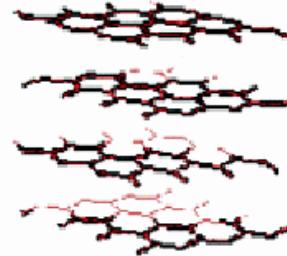
碳的一家

sp^3



金刚石

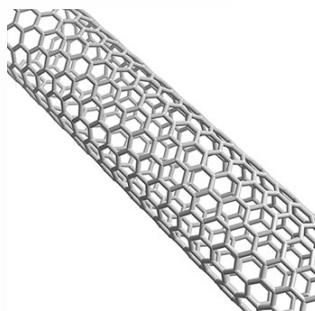
3D



石墨

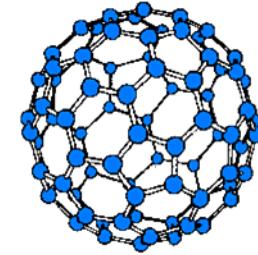
2D

sp^2



碳纳米管

1D



C₆₀

0D

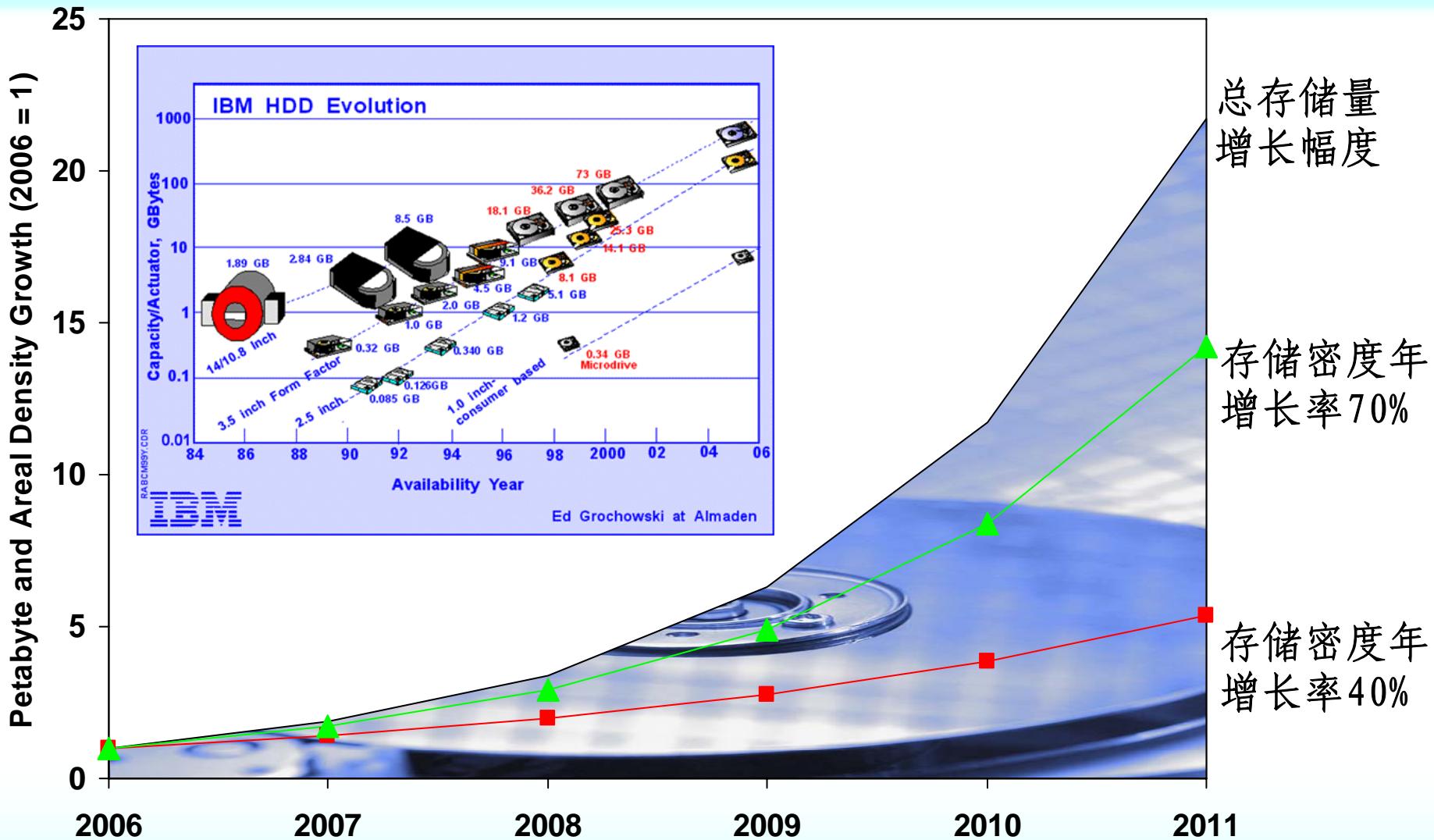
为什么对这些材料的物性感兴趣？

- 丰度：在自然界广泛存在
- 物理：是低维物理研究的理想对象，具有丰富多彩的物性
- 器件：可能用于构建更小、更快、更好的电子学器件

器件

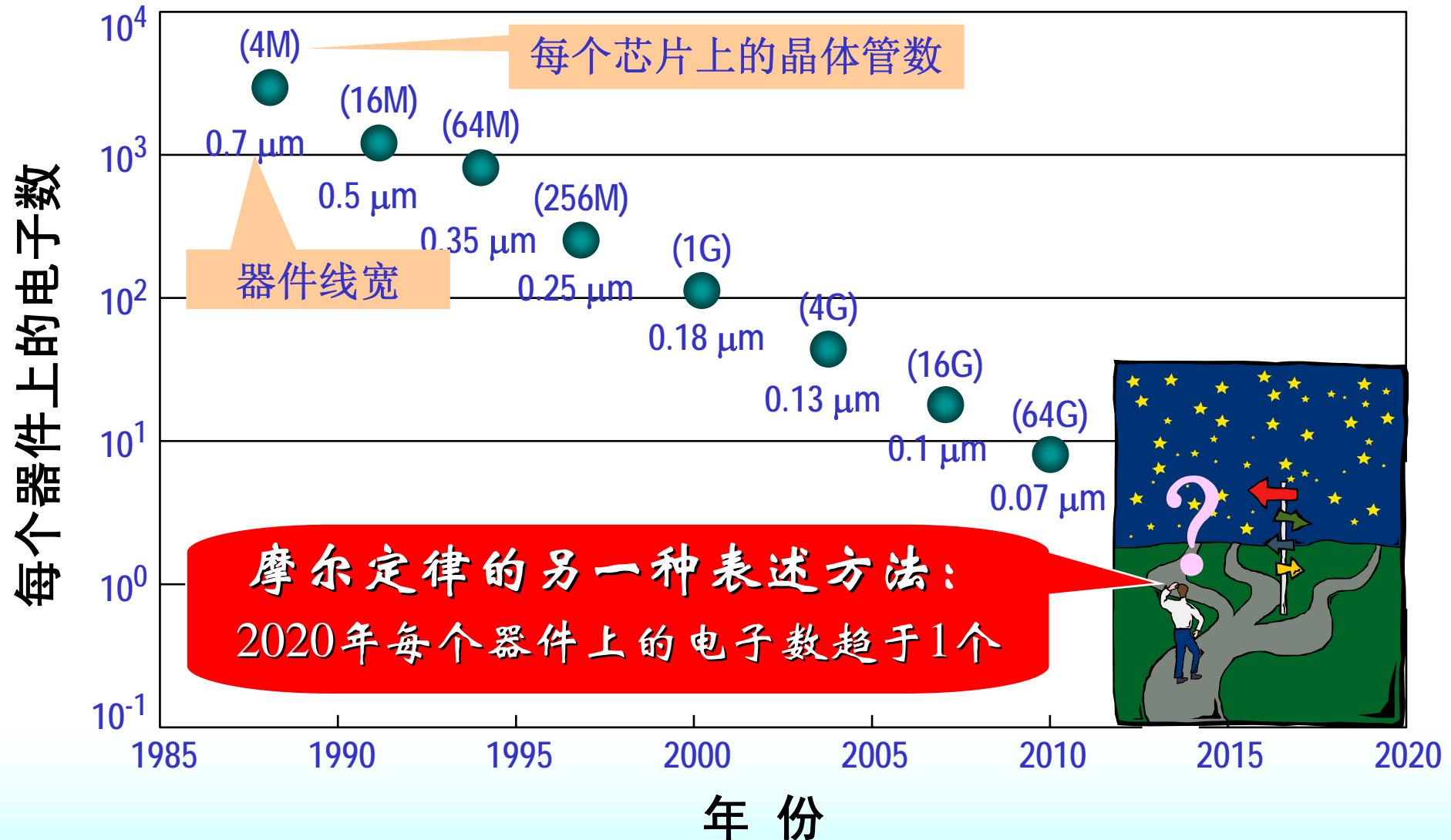
- 目前的信息技术所面临的困境
- 量子信息技术带来的希望
- 石墨烯能否用来做自旋量子比特？

五年内世界总信息存储需求量骤增20倍



Average Petabyte Growth Year on Year From 1995 to 2005 was > 85%, Seagate Analyst & Investor Meeting June 2006

器件越来越小，摩尔定律趋于极限！



解决方案

从物理学基本原理出发，解决下一代信息技术中的核心科学问题

信息载体的物理对应：

- 电子（电荷、自旋）
- 核
- 原子
- 各种人造原子（人工量子结构）
- 各种量子激元（光子， 激子， NVC, ...）

经典/半经典/量子

因为信息需要物理载体，所以

突破的希望寄托于物理学基础研究

经典/半经典/量子

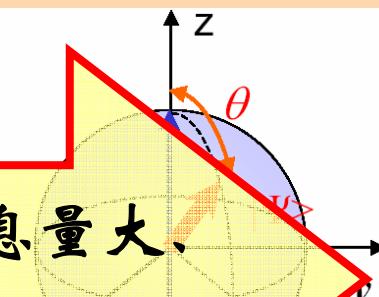
信息载体的取值空间：

量子信息技术自然具有信息量大、

不可克隆等特点，因而是重要发展方向

二进制

多进制

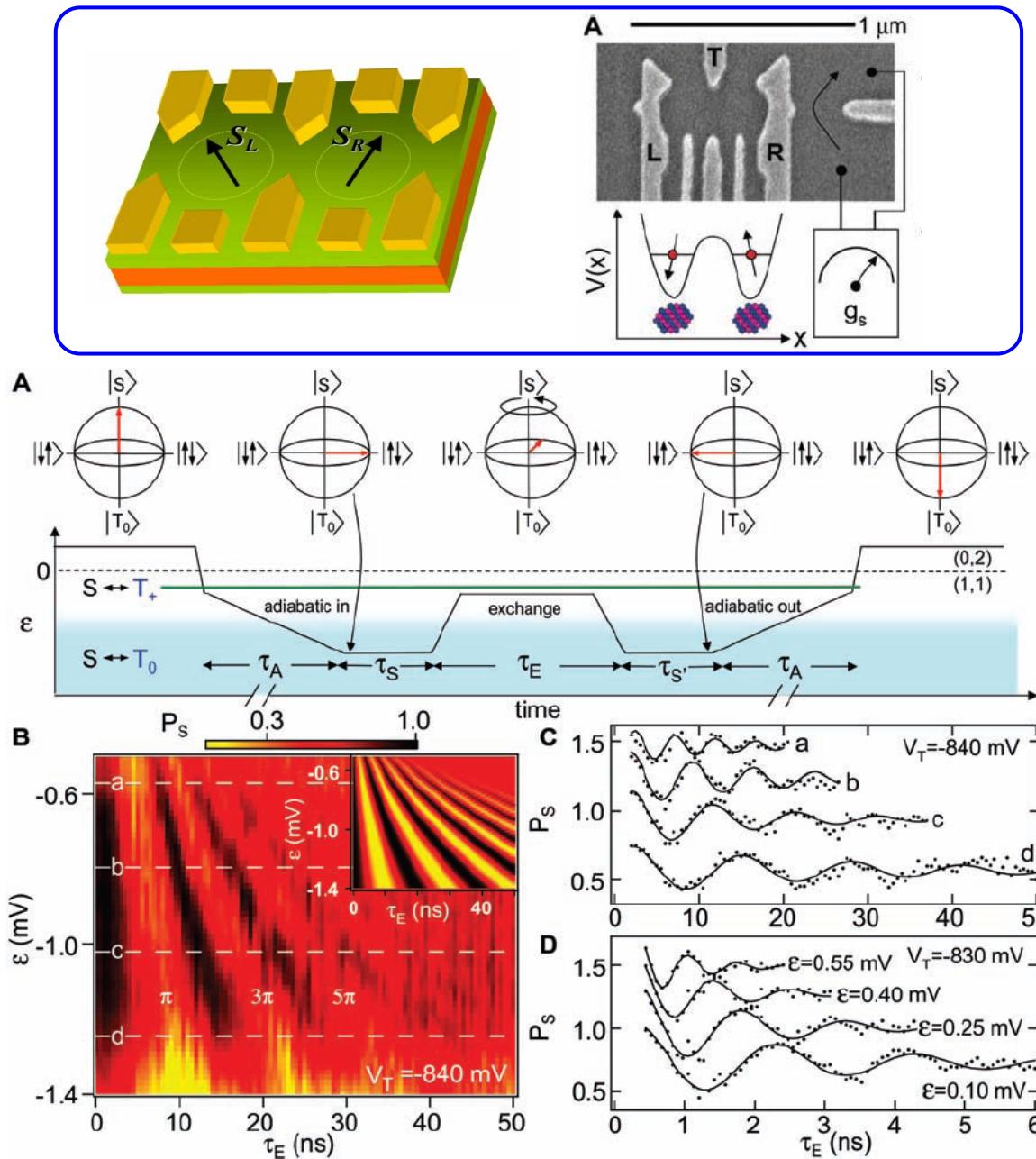


Bloch球
量子平行计算

$|\Psi|$

$|\Psi|e^{i\theta}$

Spin qubits in semiconducting quantum dots systems



AlGaAs/GaAs:

Very short decoherence time, due to:

- Hyperfine interaction
- Spin-orbital interaction

Other materials ?

- Si or Si/Ge system
- Graphene
- Carbon nanotubes

- Marcus group
- Vandersypen group
-

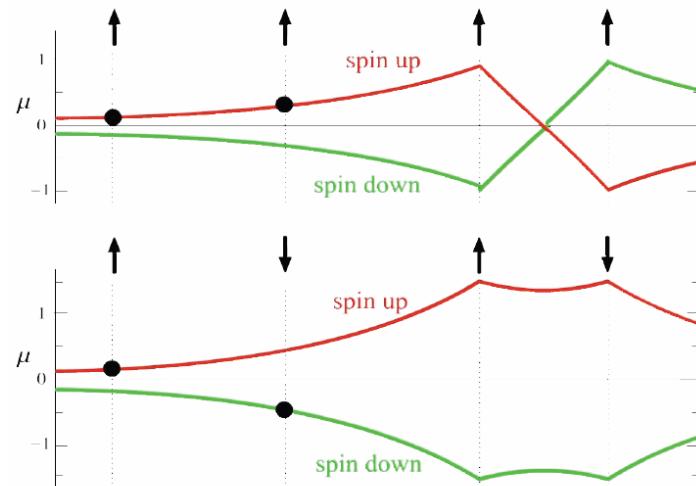
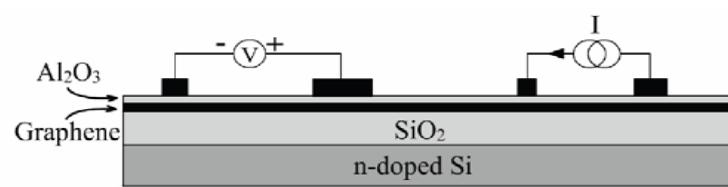
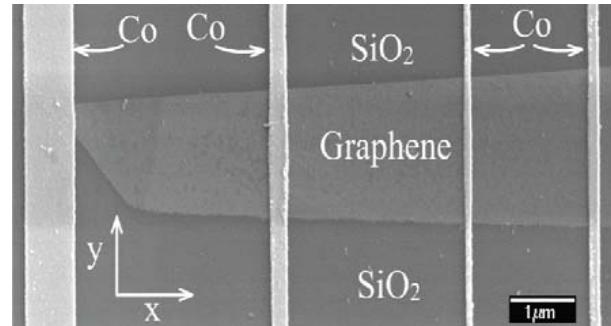
For Graphene:

- SO coupling is weak :
 $2\Delta_{\text{so}} \sim 1 \text{ } \mu\text{eV}$ (10 mK) (Y.G. Yao, et al, PRB'06)
- Hyperfine interaction is weak:
 ^{13}C (spin $\frac{1}{2}$) abundance $\sim 1\%$

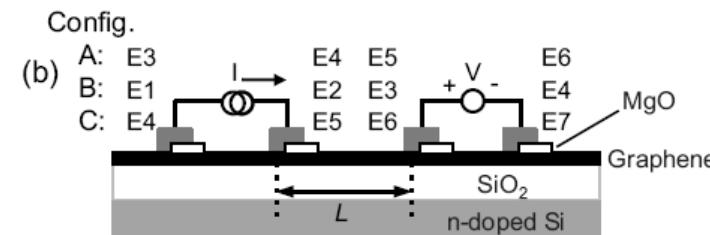
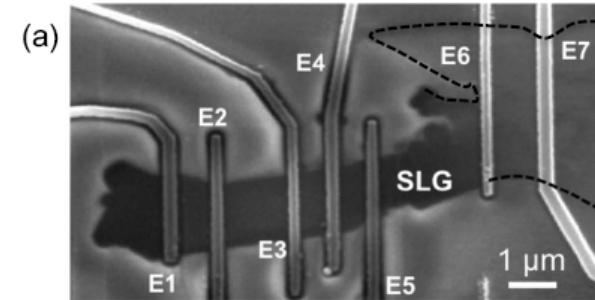
Longer T_1 and T_2 ?

Answer: No.

Spin polarized electron injection and non-local detection



van Wees group



C. N. Lau group

	van Wees	C. N. Lau
τ_{SO}	~ 150 ps	~ 84 ps
I_{SO}	~ 1.6-2.0 μm	~ 1.5 μm

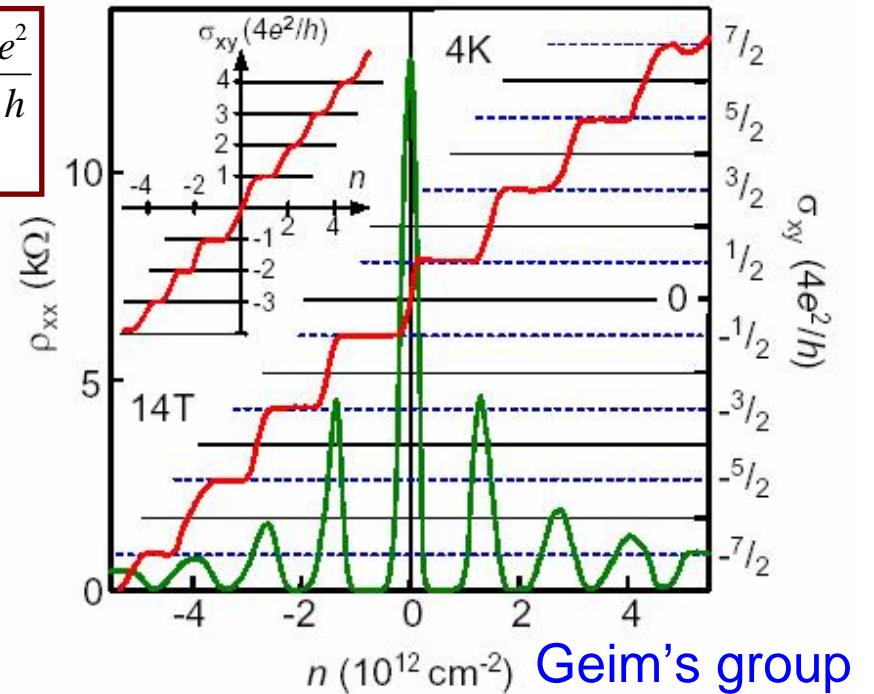
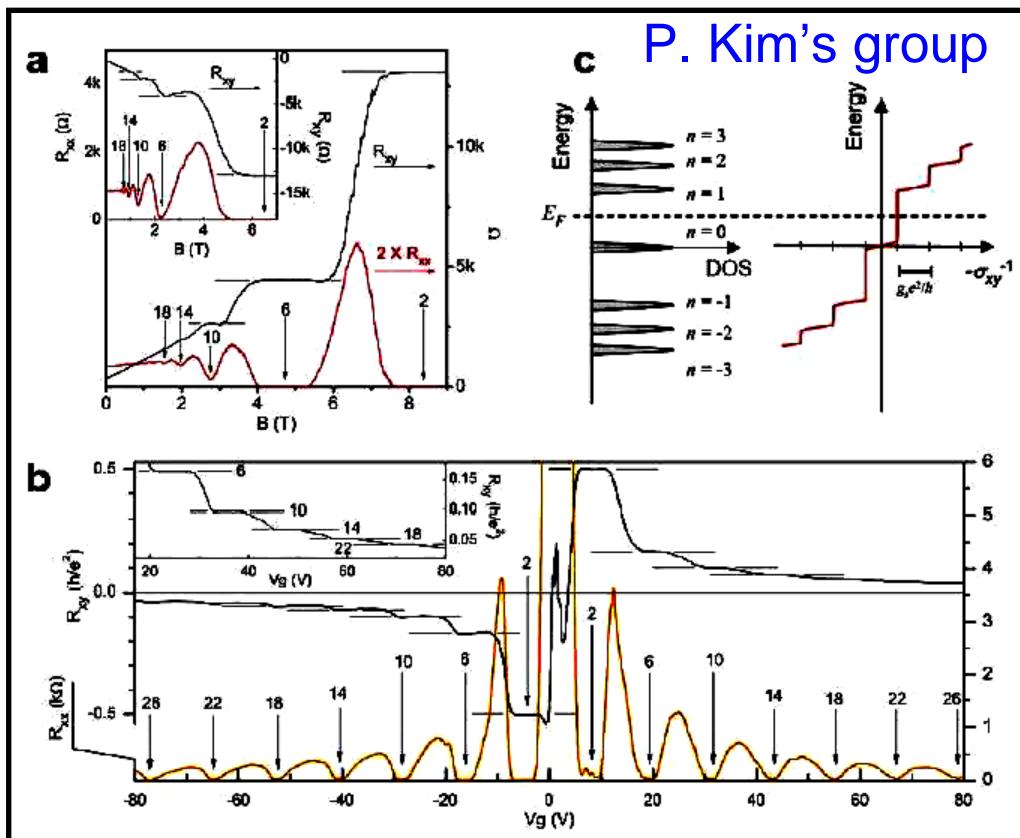
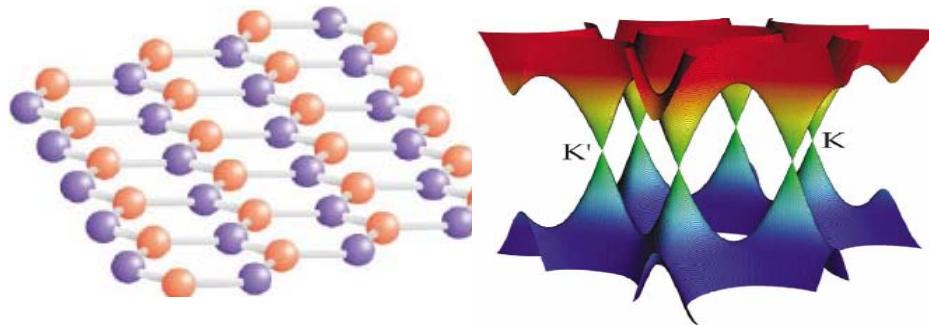
What is the cause of spin decoherence
in graphene ?

Graphene

Novel Properties:

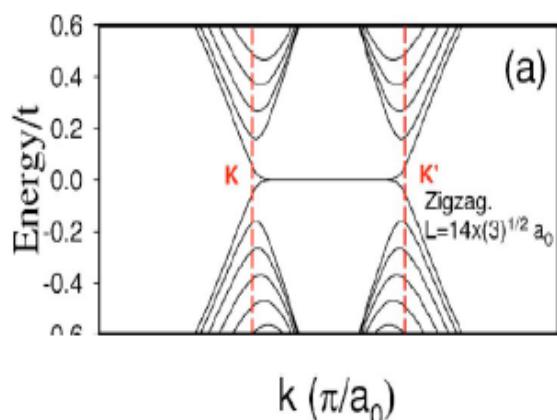
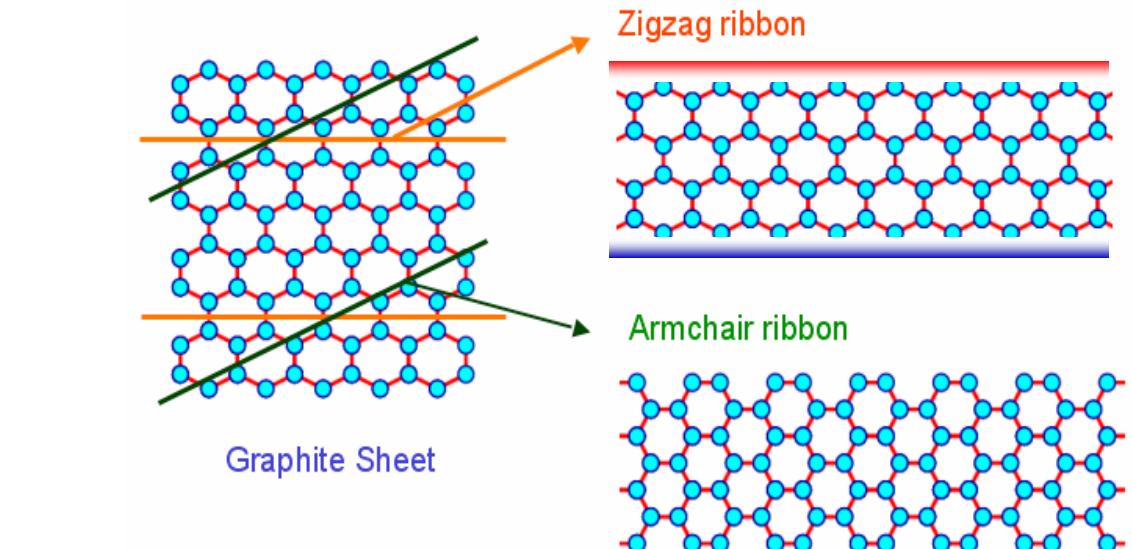
$$R_{xy}^{-1} = g_s \left(N + \frac{1}{2} \right) \frac{e^2}{h}$$

$g_s = 4$



- Two sublattices
- Pseudospin
- Four-fold degeneracy (valley x 2 + spin x 2)
- Berry phase
- Half integer QHE
- Klein Tunneling
-

Edges of graphene nanoribbon (GNR)

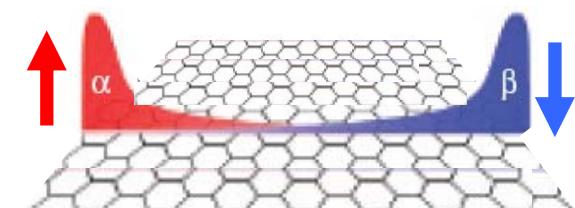
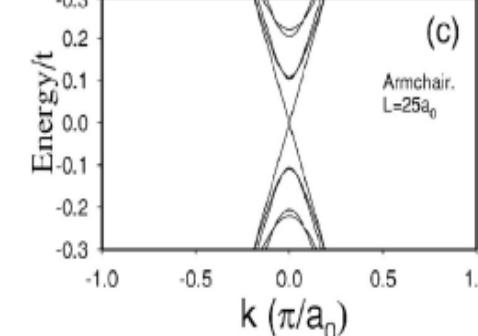
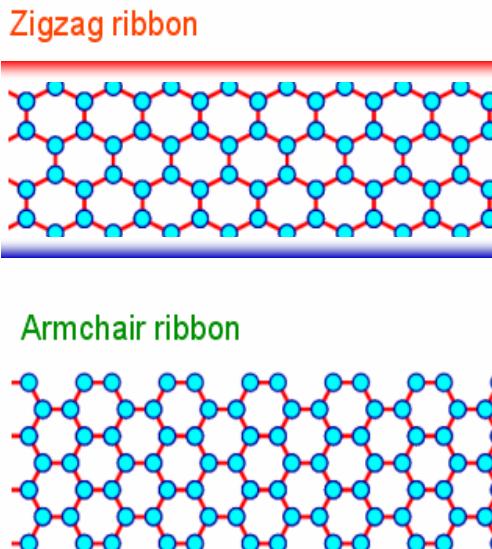


- zig-zag are gapless (edge modes)
- armchairs: gapped or gapless

$W = 3N, 3N+2$

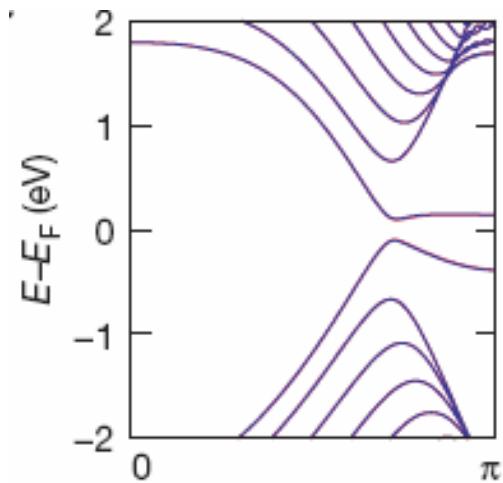
$W = 3N+1$

Ezawa, PRB 2006; Brey-Fertig, PRB 2006

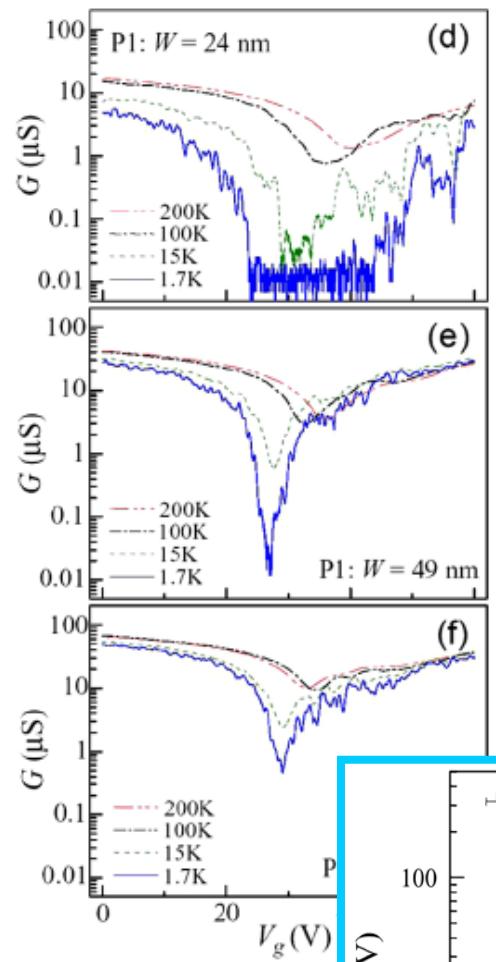
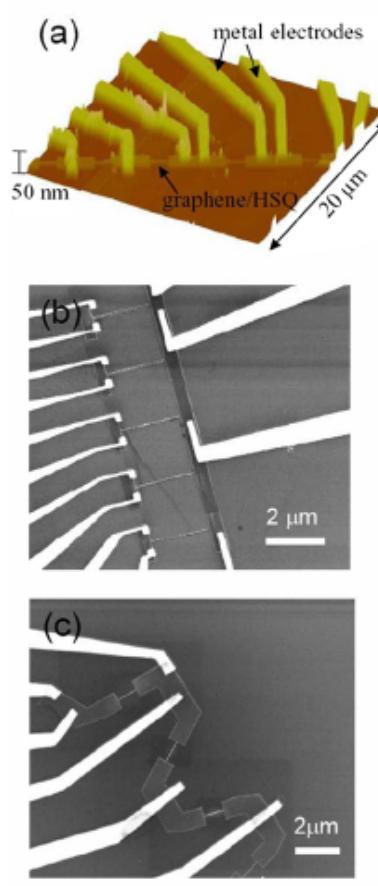


Theoretical prediction:

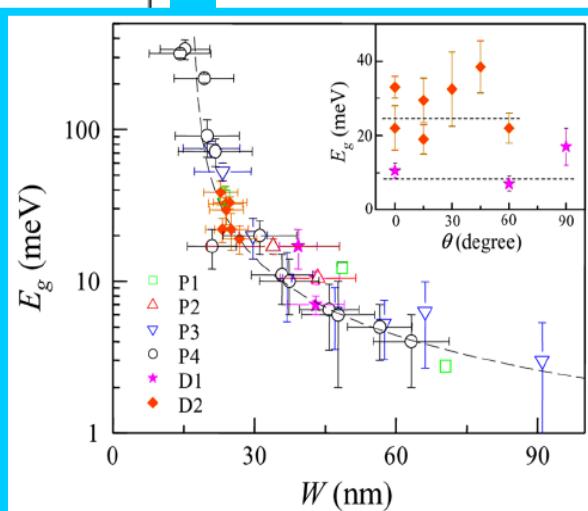
- Zigzag GNR has spin polarized edge states.
W.Y. Son et al., Nature'06
L. Pisani et al., PRB'07
- Same for bilayer GNR.
E.V. Castro et al., PRL'08



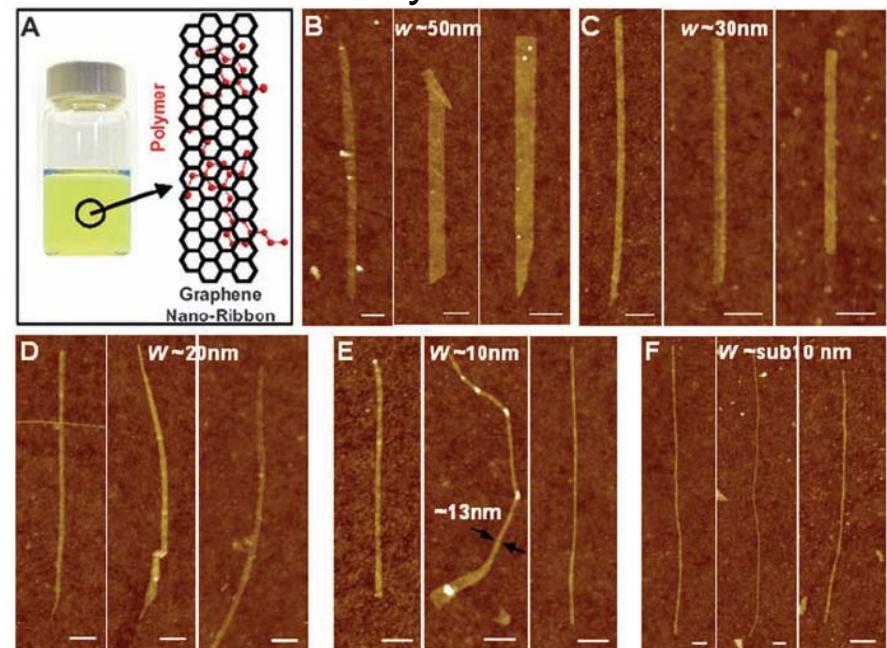
Experiments on GNR



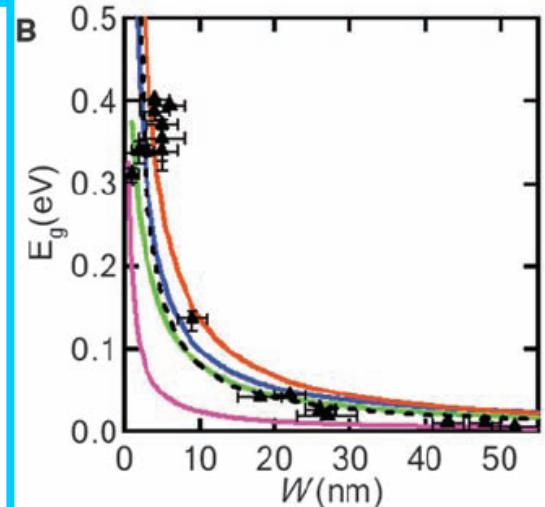
Kim Group, PRL'07
Lithograph-made GNRs



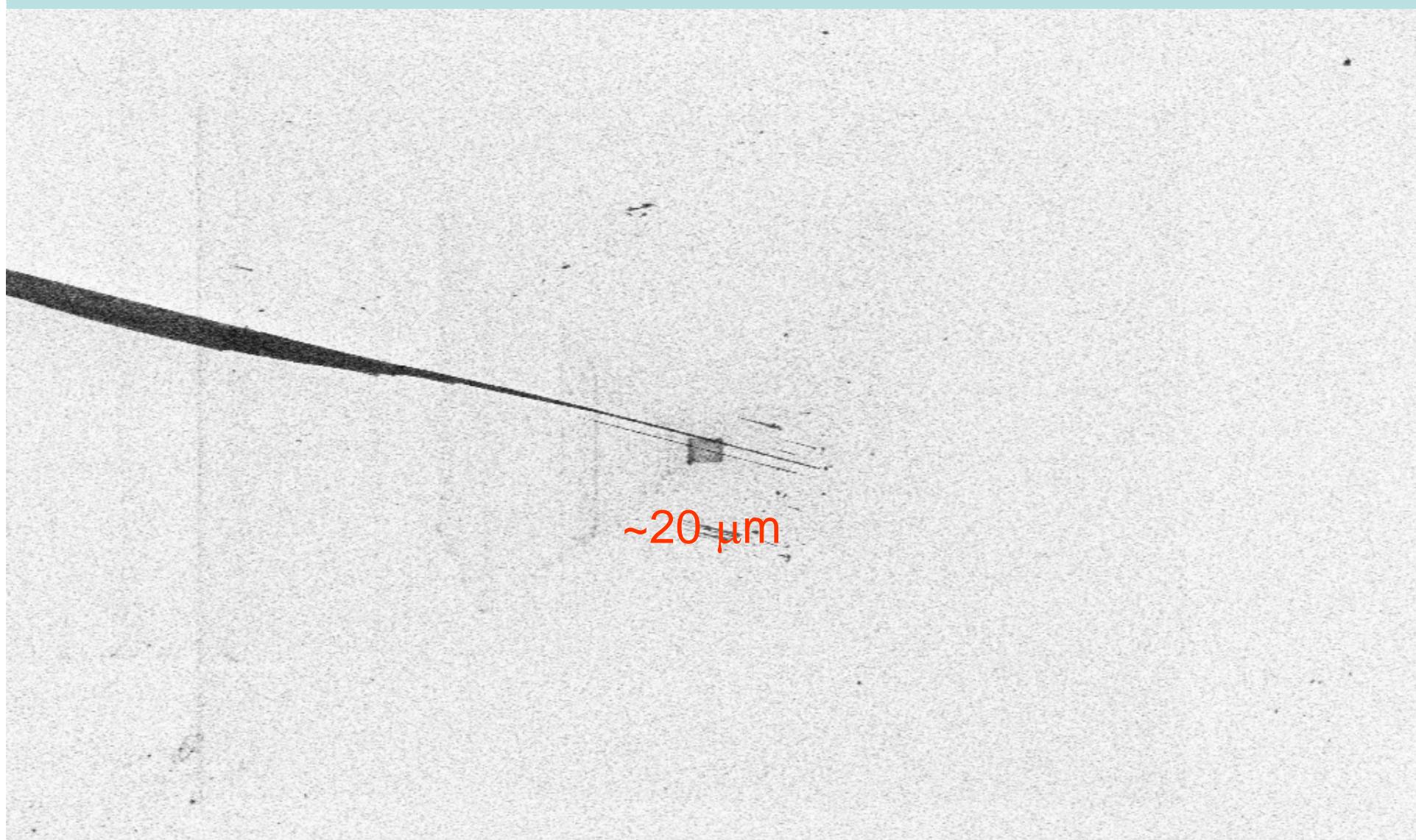
Chemically derived GNRs



Dai Group, Science'08



Our samples: mechanically exfoliated GNRs



~20 μm

Raith 150
Mag = 713 X

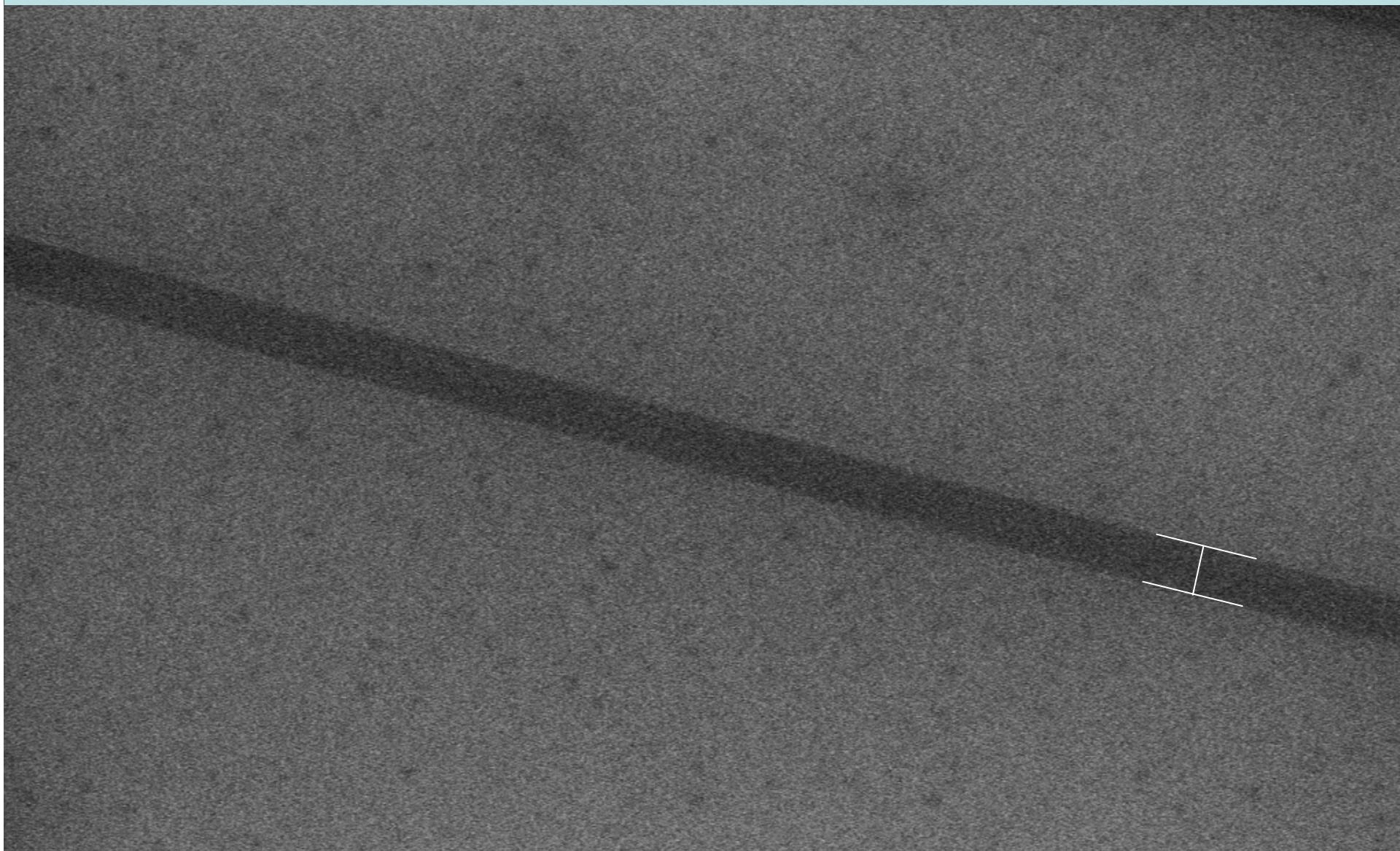
10 μm

EHT = 5.00 kV
WD = 5 mm

Signal A = InLens
User Name = TRAINING

Date :16 Oct 2008
Time :9:58:10

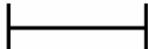
Our samples: mechanically exfoliated GNRs



Raith 150

Mag = 42.74 K X

200nm



EHT = 5.00 kV

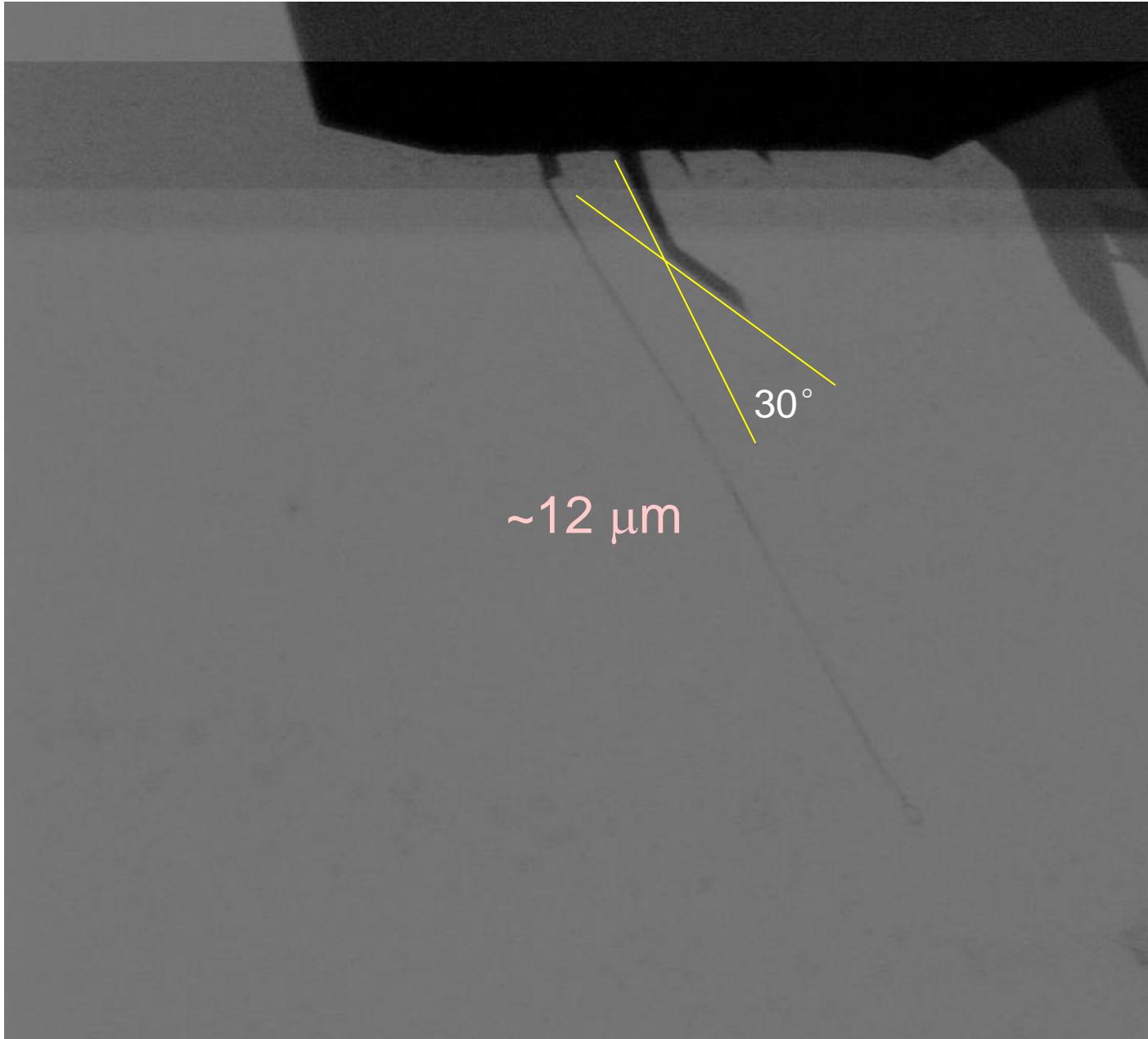
WD = 5 mm

Signal A = InLens

User Name = TRAINING

Date : 16 Oct 2008

Time : 9:56:36



The GNR used in this experiment:

- With atomic-level smooth edges
- Along principle axes

$\sim 12 \mu\text{m}$

Raith 150

Mag = 4.10 K X

2 μm^*



EHT = 5.00 kV

WD = 5 mm

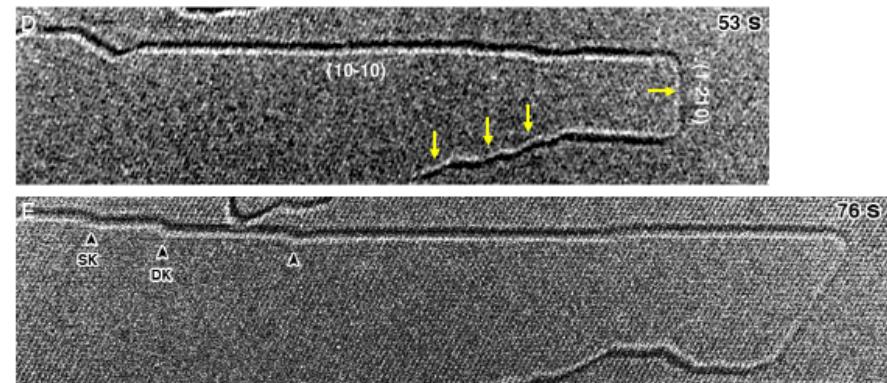
Signal A = InLens

User Name = TRAINING

Date : 11 Jan 2008

Time : 10:05:37

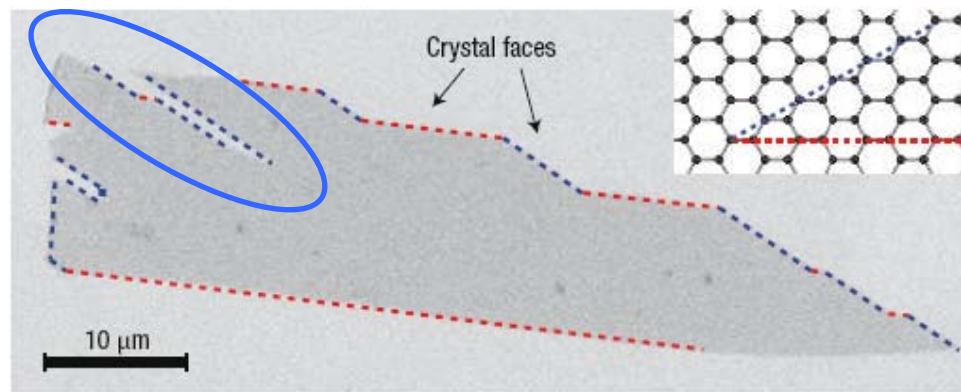
Edges of graphene nanoribbon (GNR)



99% edges are zigzag

Jian Yu Huang et al., PNAS'09

Zigzag ribbon

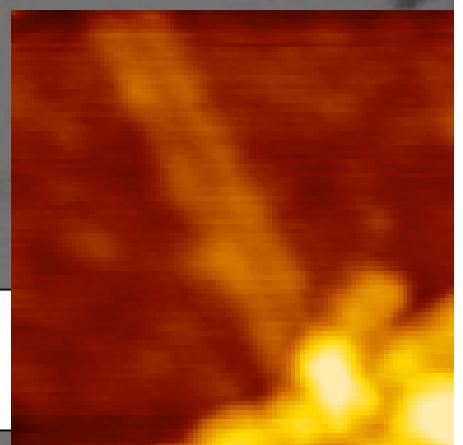
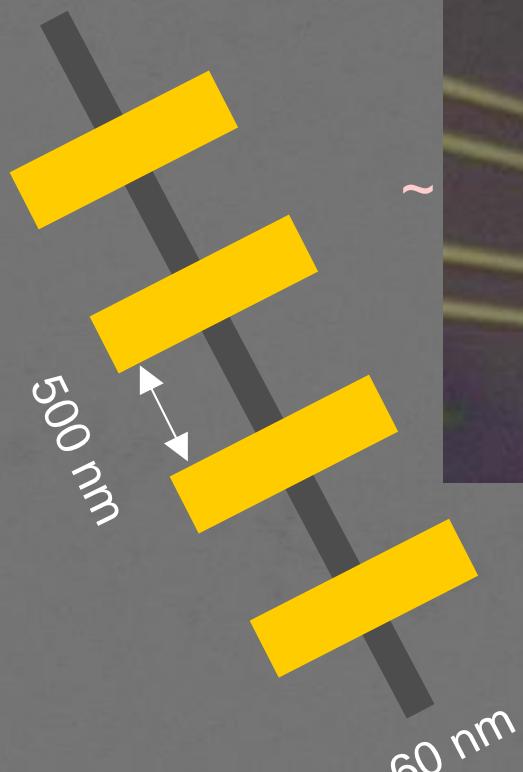


Edges are along principle axes.

A. K. Geim, Nature Materials'07

The GNR used in this experiment:

- With atomic-level smooth edges
- Along principle axes



Raith 150

Mag = 4.10 K X

2 μ m*

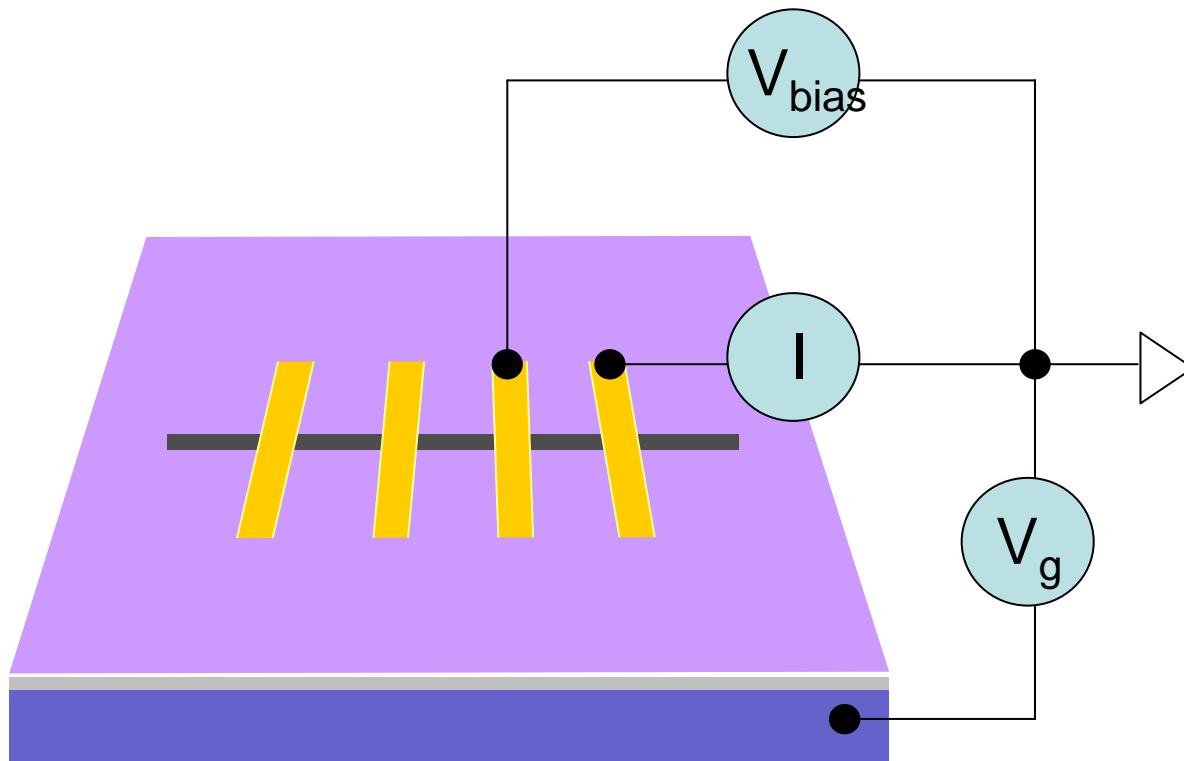


EHT = 5.00 kV

WD = 5 mm

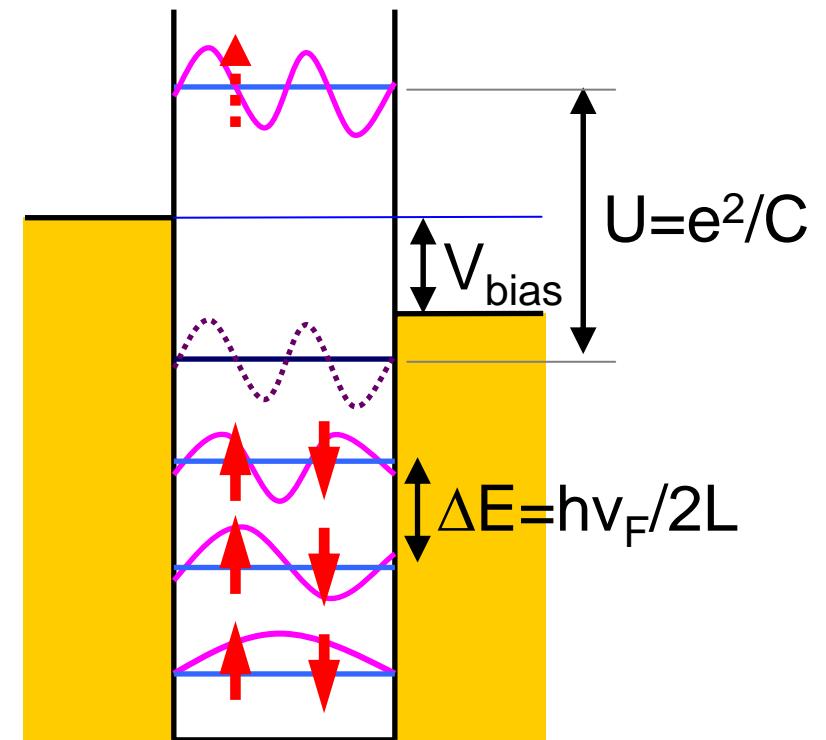
Signal A = InLens

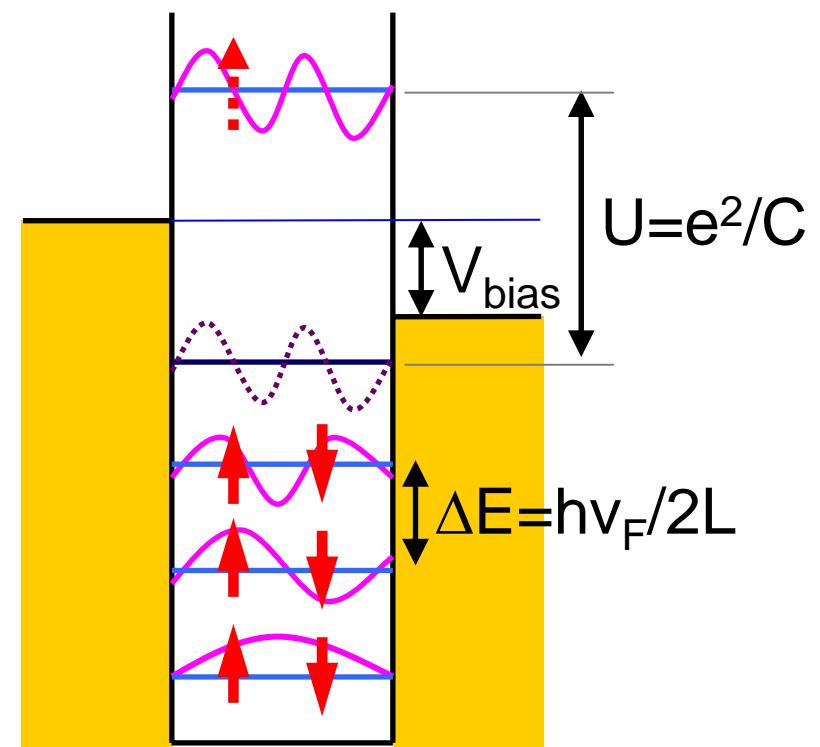
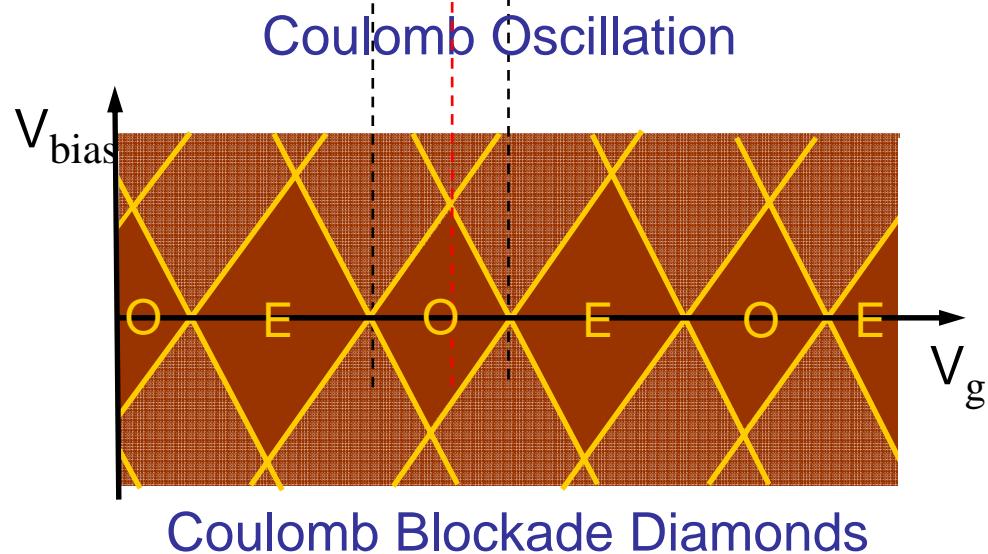
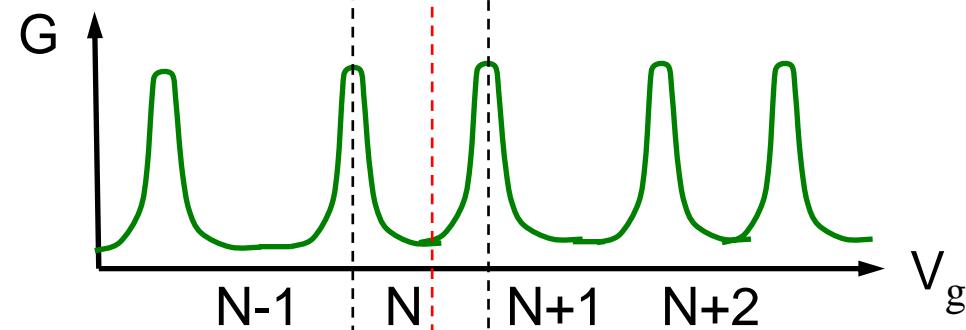
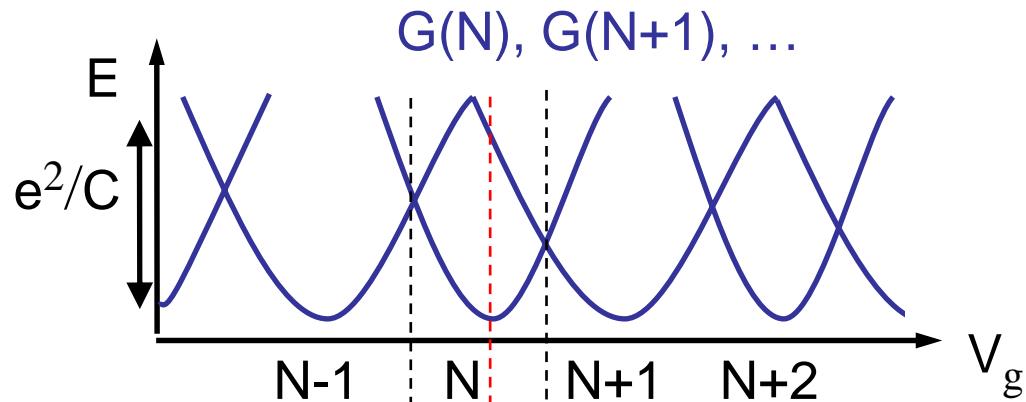
User Name = TRAINING



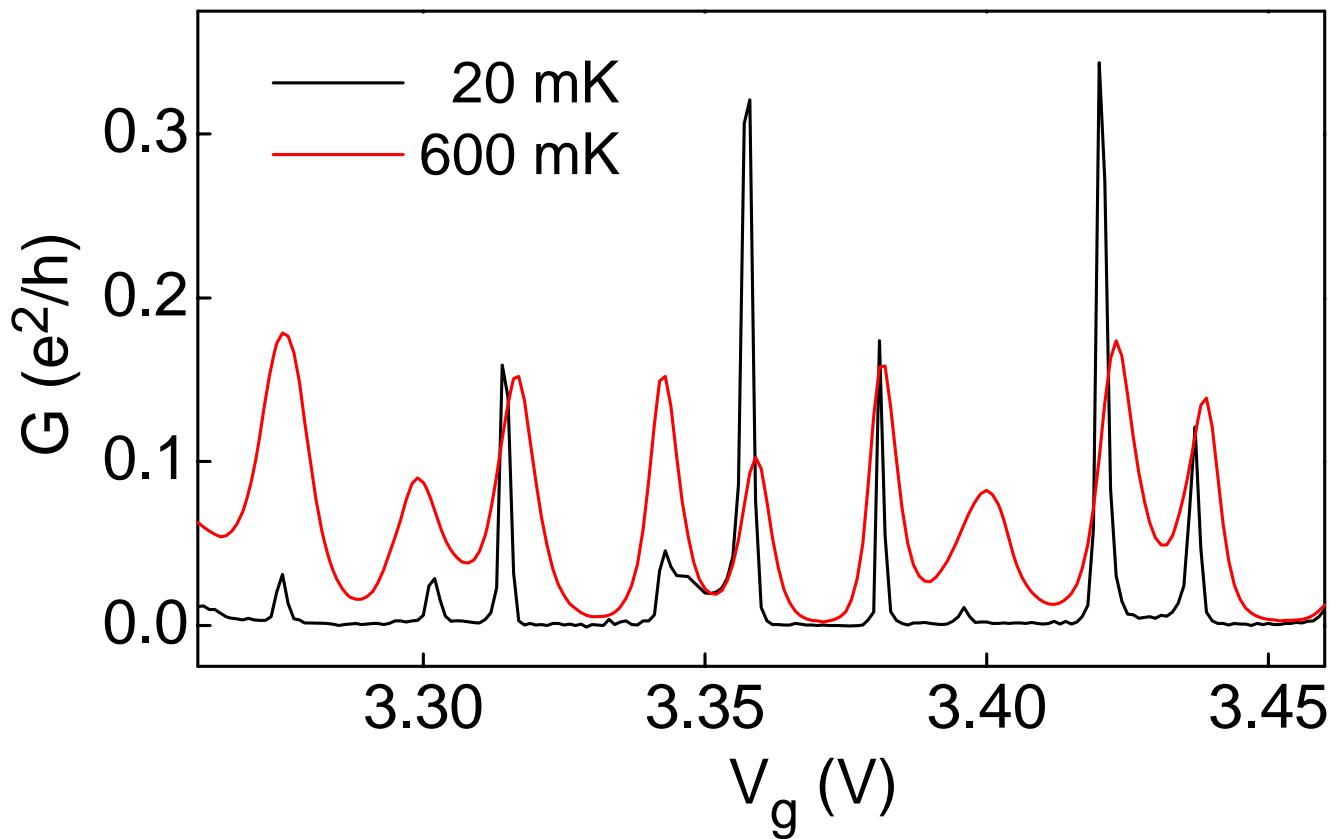
Interplay of different energy scales:

- ΔE , level spacing (excited energy)
- V_g , shifts the levels in the QD
- V_{bias} , the bias window
- U , Coulomb charging energy

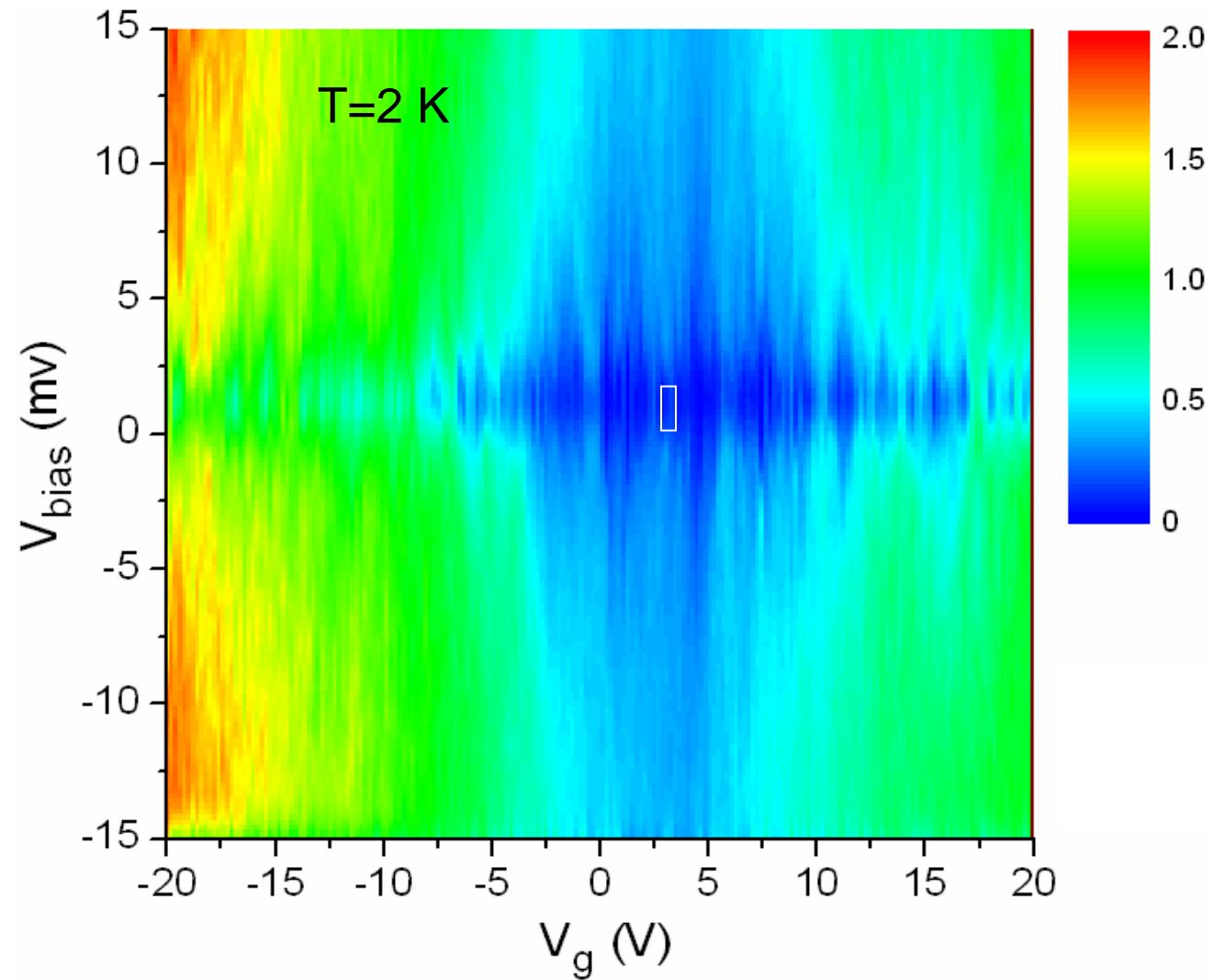




Coulomb oscillation: G peak gets sharper at low T

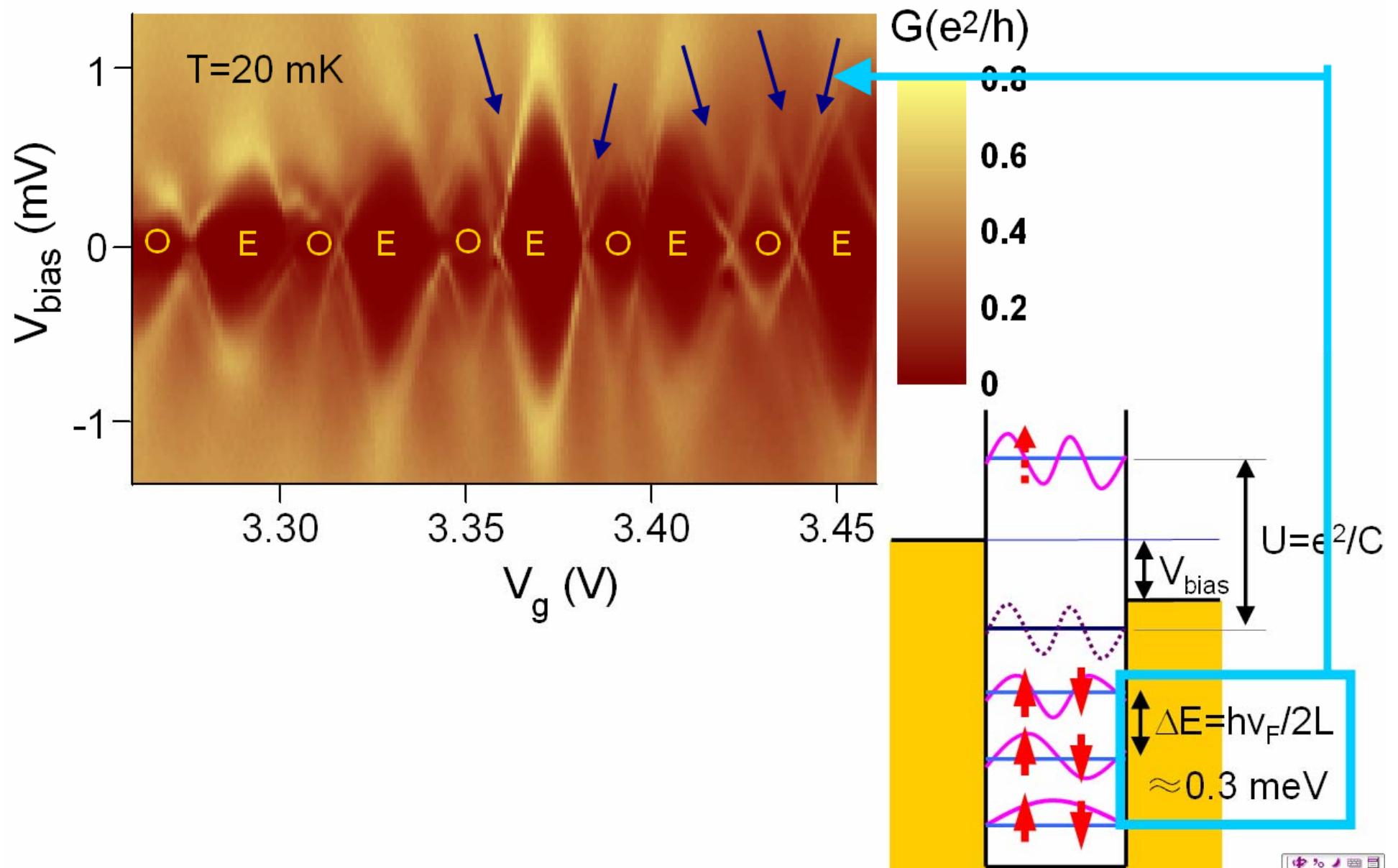


Conductance as a function of V_{bias} and V_g

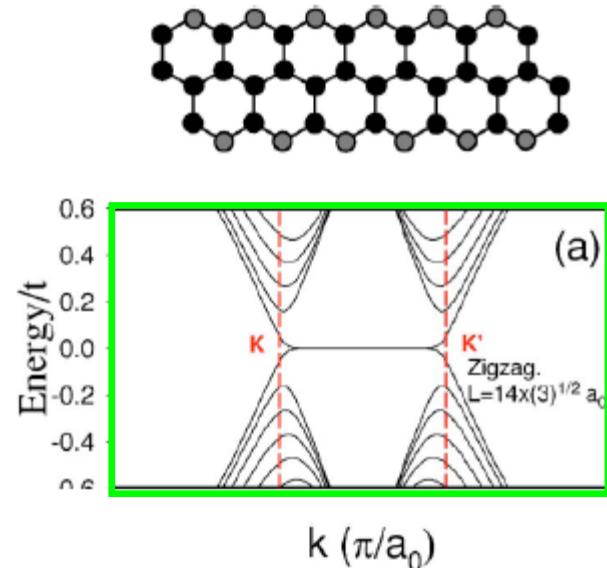


Two-fold spin shell filling in a bar-like QD

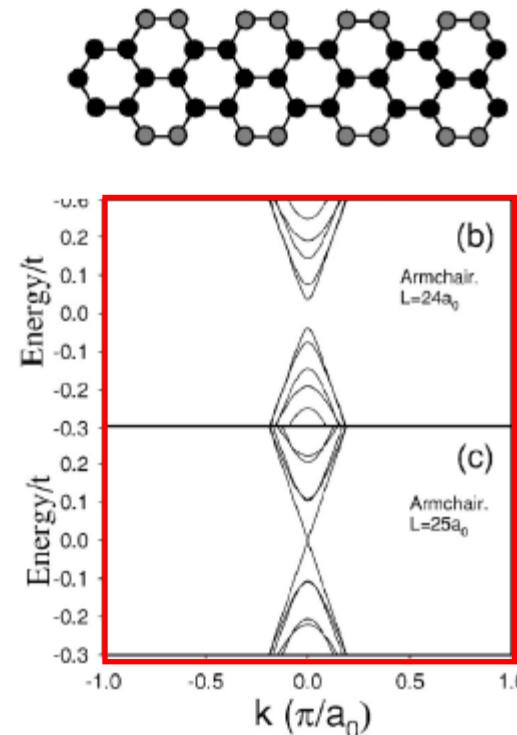
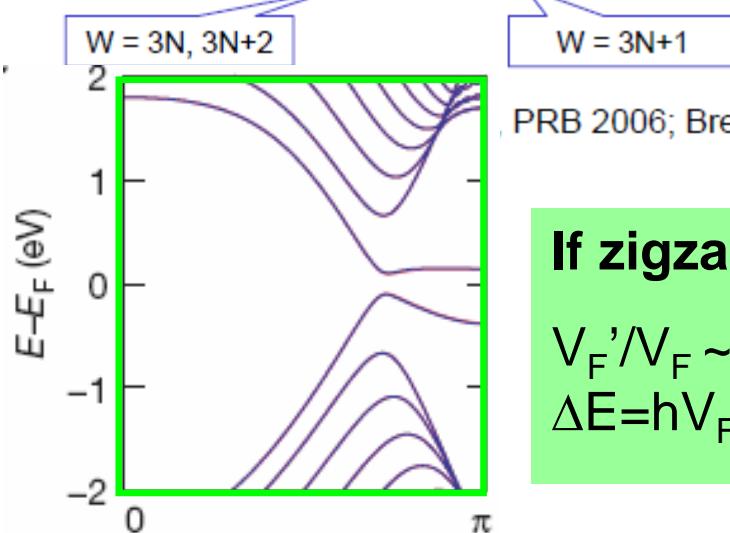
→ Break down of K and K' degeneracy



Excited energy → Zigzag ribbon



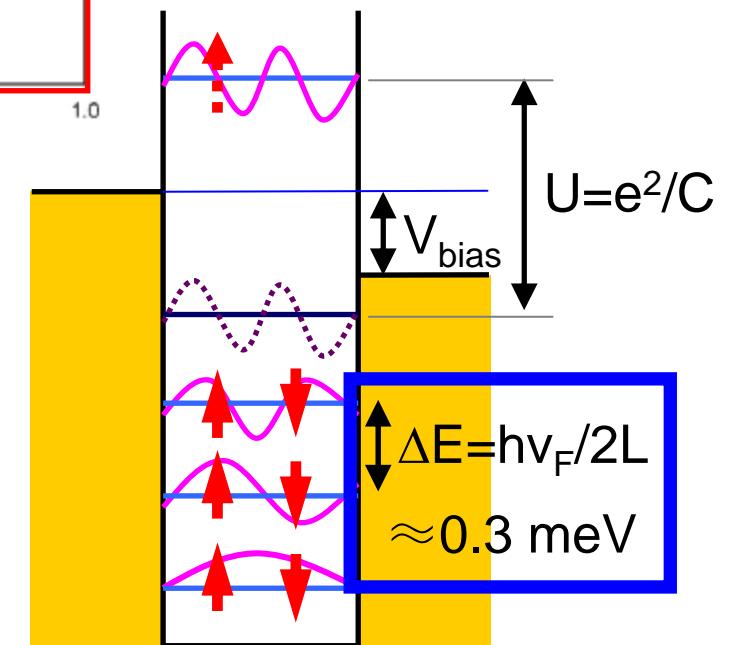
- zig-zag are gapless (edge modes)
- armchairs: gapped or gapless



If armchair:

$$V_F = 1 \text{e}6 \text{m/s}, \Delta E \approx 12 \text{meV}$$

To get $\Delta E \approx 0.3 \text{meV}$, more than 10 subbands must be filled with ~ 700 elec.
i.e., Dirac point is 10V away from the V_g window.

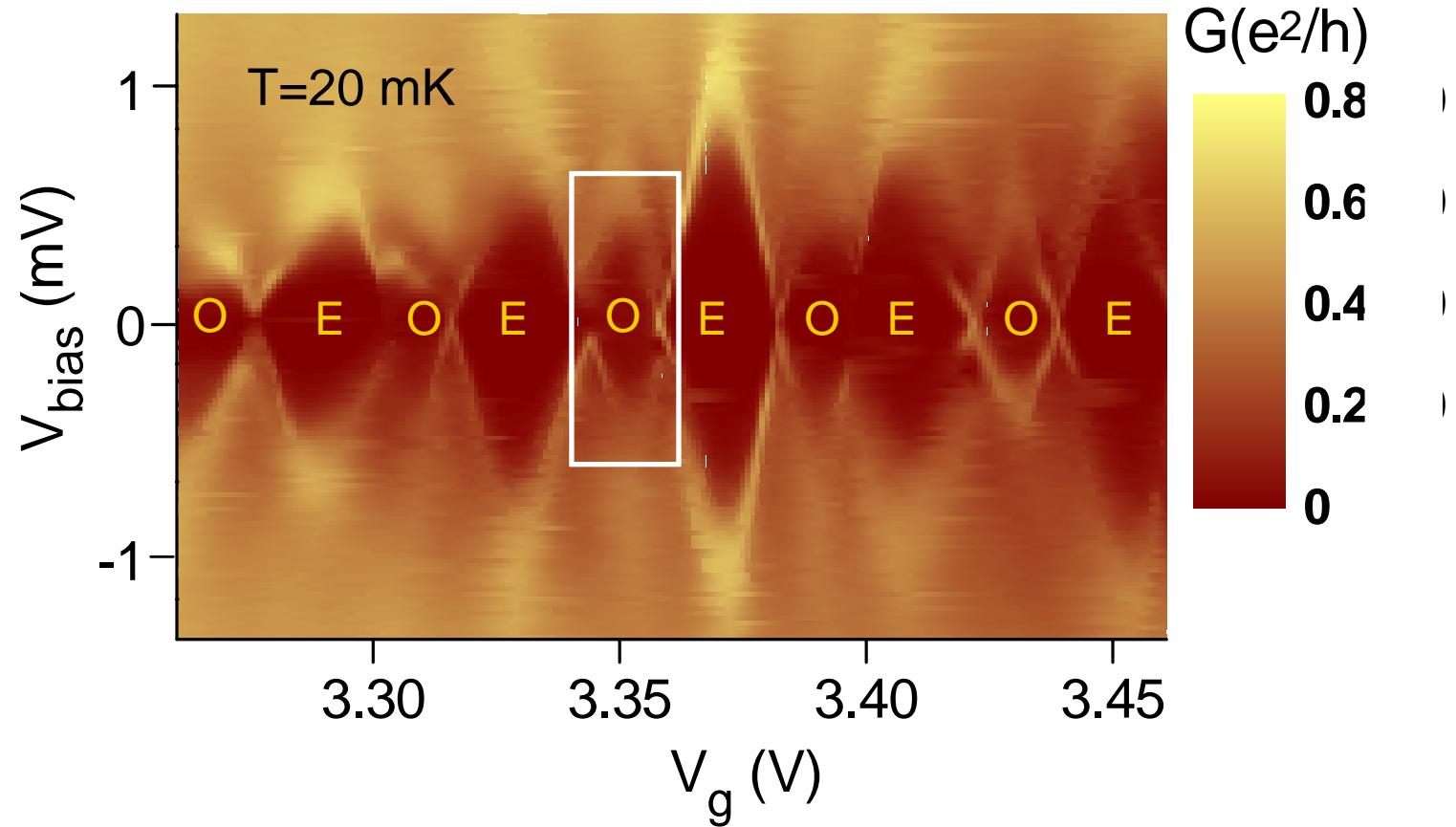


If zigzag:

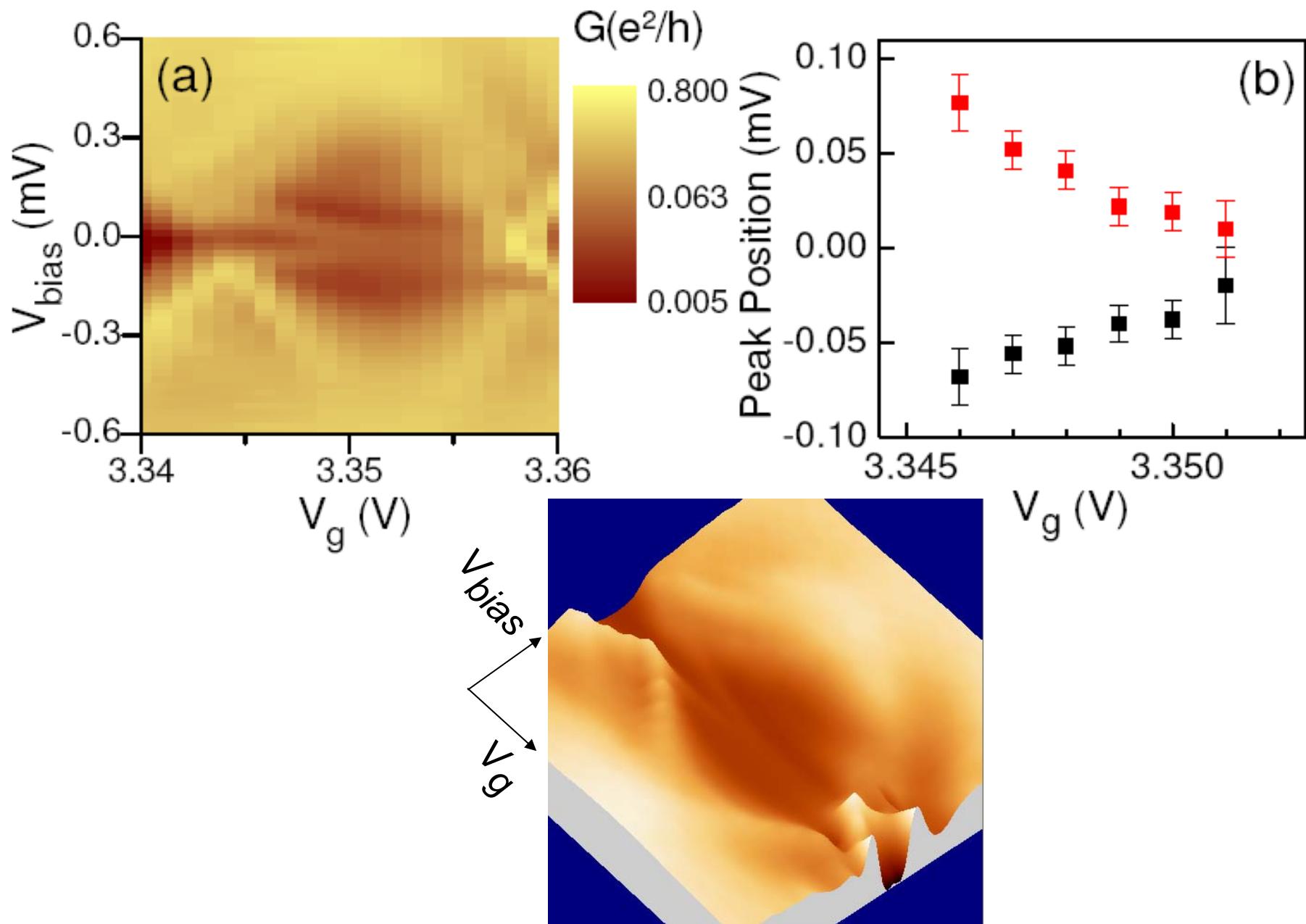
$$V_F'/V_F \sim t'/t \approx 1/30$$

$$\Delta E = hV_F'/2L^* \approx 0.4 \text{ meV}$$

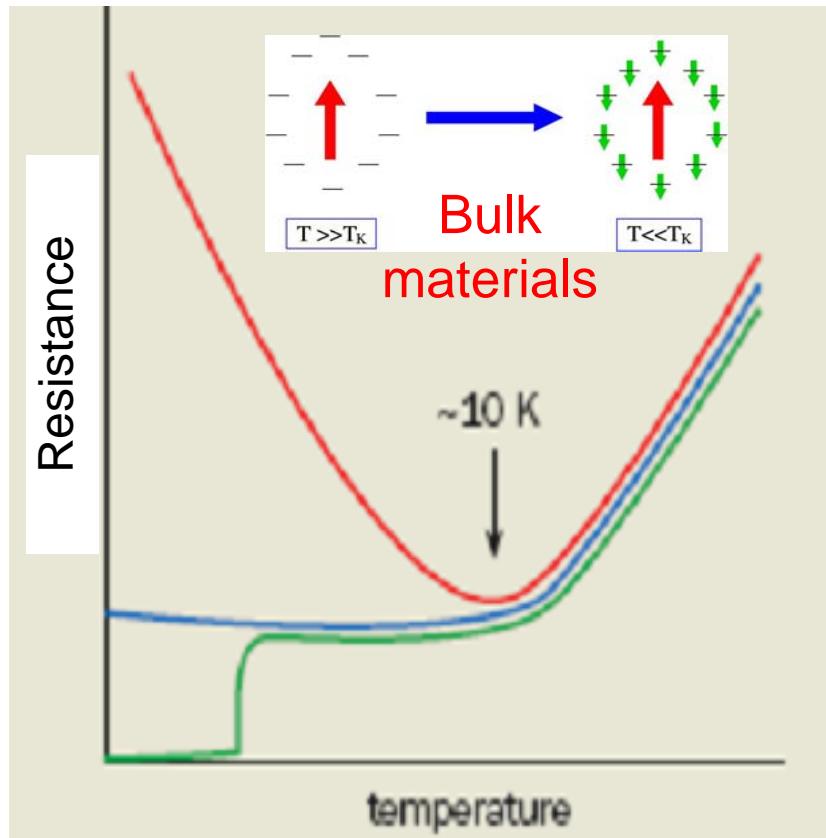
Kondo resonance



Inelastic Kondo resonance



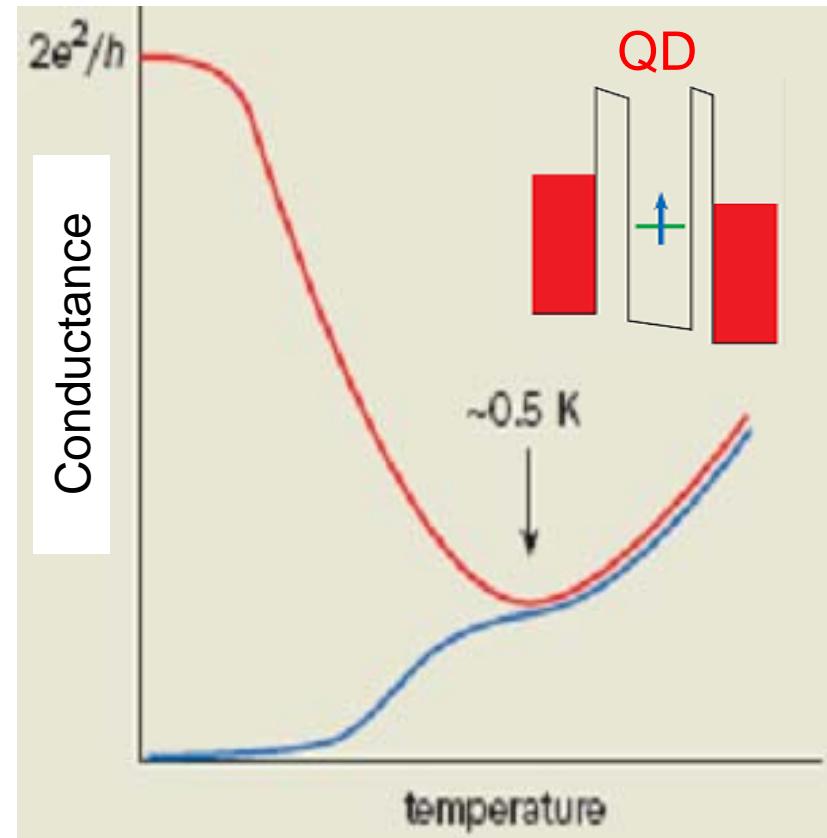
Kondo problem



Scattering and screening:

$R \uparrow @ \text{low } T$

Kondo

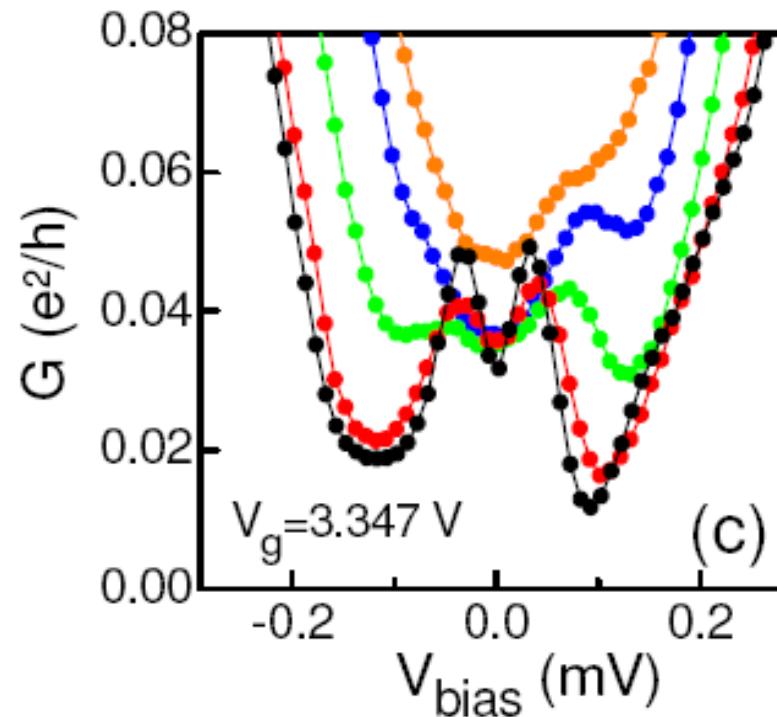


Resonant co-tunneling:

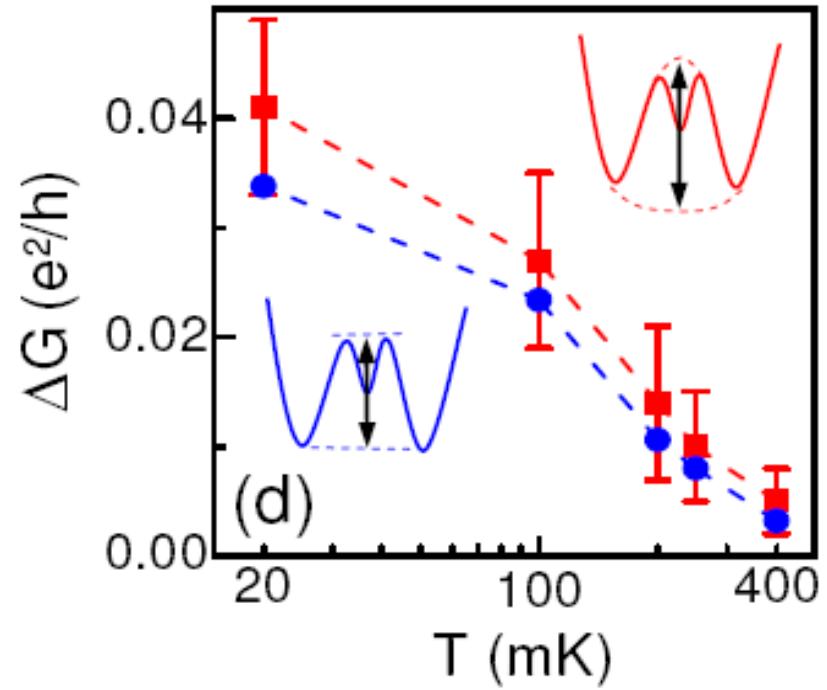
Provides a conduction mechanism in the CB regime.

T. K. Ng and P. A. Lee' 1988
L. I. Glazman, et al.'1988

Logarithmic temperature dependence

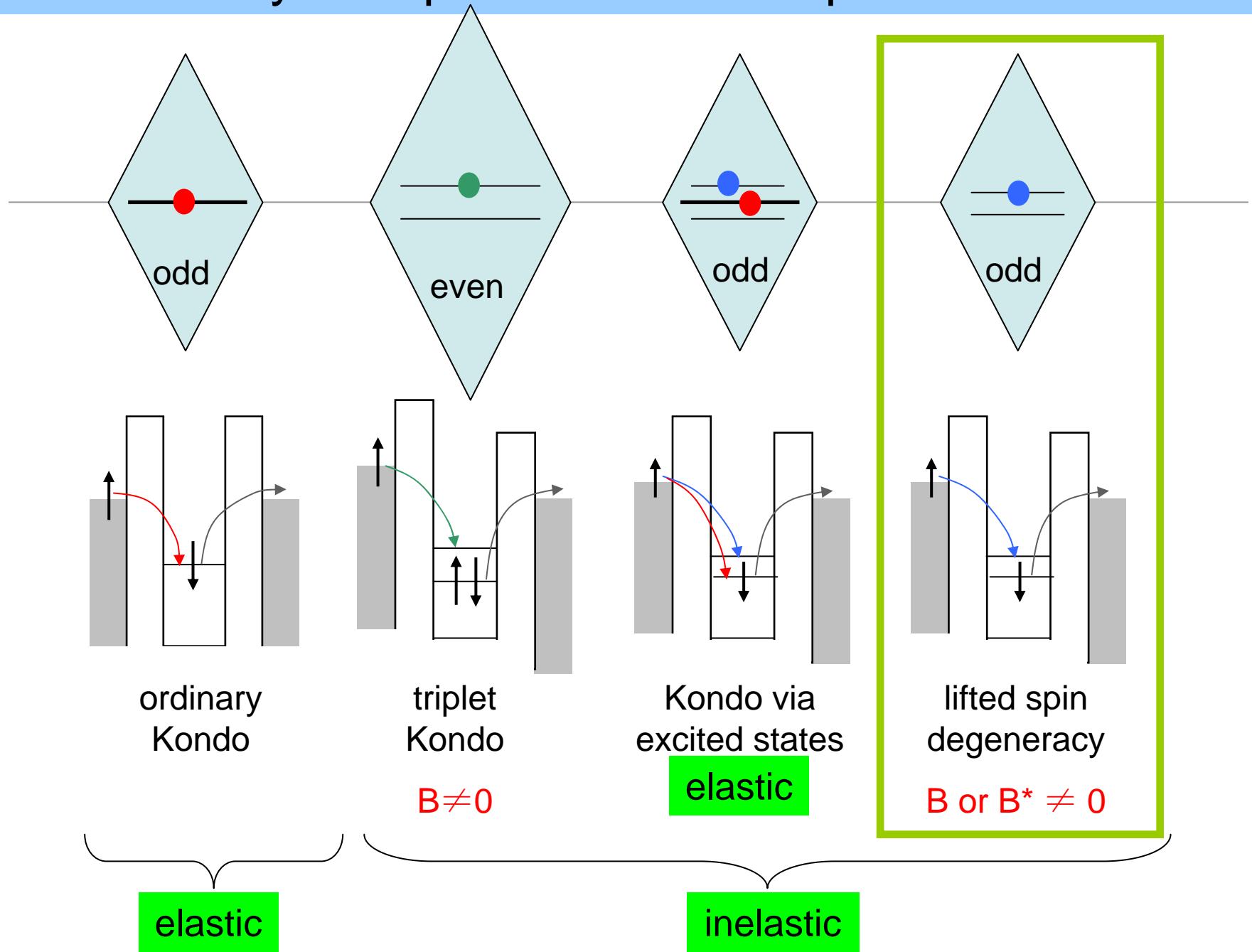


T (mK)
—●— 400
—●— 250
—●— 200
—●— 100
—●— 20

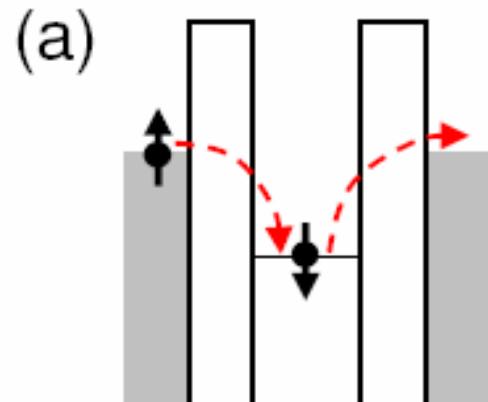


Kondo-like

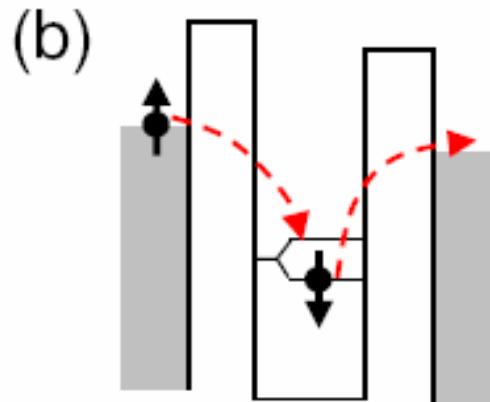
Why multiple conductance peaks ?



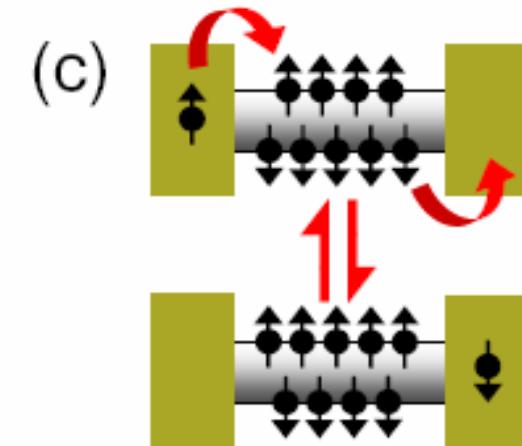
A toy model for inelastic Kondo resonance involving spin-polarized edge states



Elastic Kondo

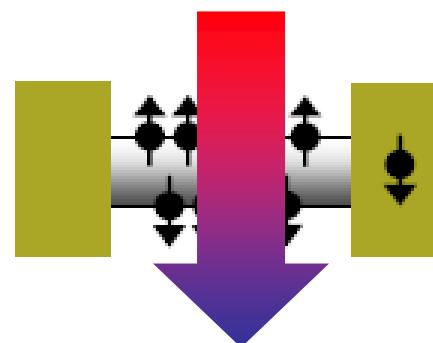


Inelastic Kondo



Inelastic Kondo realized in a zigzag GNR at $B=0$,

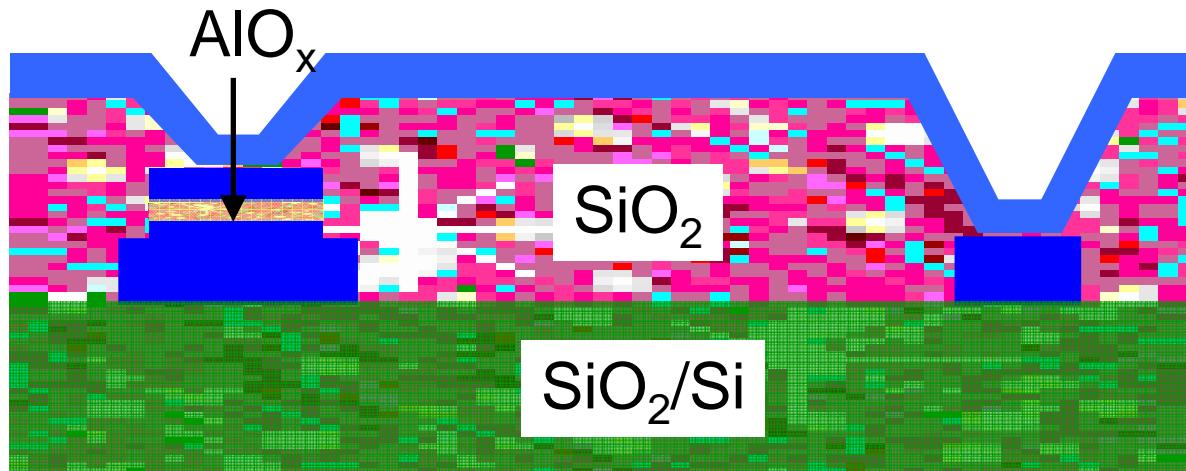
When the electrostatic potential of the two edges are different,
due to applied transverse electric field or trapped charge vacancies.



Spin qubits made of graphene quantum dots ?

- There are mobile charges in SiO_2

Decoherence
by S. Oh



- Unfortunately the spins in zigzag GNRs are coupled to mobile charges in the environment.

Summary

- Measured the transport properties of a zigzag GNR.
- Observed two-fold spin shell filling in CB regime.
- Provided transport evidence of spin polarization in the zigzag GNR.

C. L. Tan, et al, submitted

感谢



合作者： 学生：

张殿琳	周 锋
解思深	刘首鹏
杨昌黎	谭长玲
顾长志	谭振兵
	马 丽
	杨 帆
	王 轲
	刘 恒